

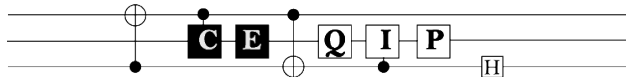
# Geometry in Entanglement Percolation

John Lapeyre

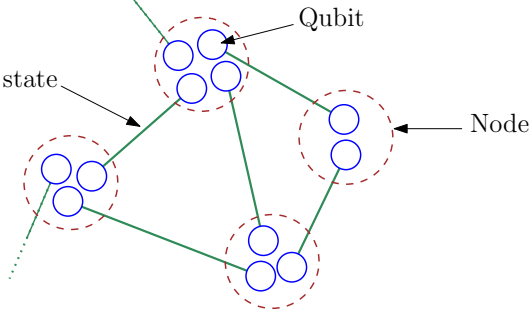
ICFO, Barcelona, Spain

June 8, 2014

CEQIP, Znojmo, Czech Republic



Bi-partite entangled state



- *qubits* (two-level quantum systems) (spin  $1/2$ , e.g. photon polarization)
- Multiple qubits and classical resources at each *node* (vertex)
- *links* (edges): bi-partite entangled (pure/mixed) two-qubit states.
- Goal: entangle pairs of qubits between distant nodes
- Quantum operations local: within nodes. Classical can be global.  
Local Operations and Classical Communication **LOCC**

## Technical motivation: Generalize one-dimensional networks

- Quantum Information: Entanglement is a resource for tasks: teleportation, key distribution, fault tolerant computation
  - Creating entanglement requires local interaction. Noise increases with distance. Depolarization. Absorption. **Can't distribute entanglement over long distance in a single stage!**
- Long range entanglement via Network of stations or nodes that store and purify a state.
  - Generalization of **quantum repeater** schemes.  
*Dür, Briegel, Cirac, Zoller, PRA 1999*
  - Nodes share partially entangled states of qubits
  - Nodes(stations)/channels, Vertices/edges, Sites/bonds
  - Quantum operations **probabilistic**
  - Large number of random components  $\Rightarrow$  **Complex Networks, Percolation, Phase transition**

# Goal of Entanglement Percolation

- Given a network with a specified amount of quantum and classical resources, and a specific long range entanglement task, **design the optimal protocol to achieve the task.**
- E.g. Optimal: Smallest amount of resources (entanglement) per link that achieves task. Or protocol that achieves task with highest probability for a given amount of resources.
- E.g. Topology of lattice(network) may be an external constraint.
- E.g. Task: entangle fixed widely separated nodes A and B.

# Entanglement: Two entangled qubits



Two entangled qubits: four-dimensional Hilbert space.

## Bi-partite pure state

All such states LOCC equivalent to unique state in Schmidt basis.

$$|\alpha\rangle = \sqrt{\alpha_0} |00\rangle + \sqrt{\alpha_1} |11\rangle$$

$$\alpha_0 > \alpha_1 \quad \alpha_0 + \alpha_1 = 1 \quad \alpha_1 \in [0, 1/2]$$

Pure, partially entangled, bipartite state

$\alpha_1 = 0$ : no entanglement,  $\alpha_1 = 1/2$ : max. entanglement

# Bell State: Singlet Conversion



Partially Entangled:  $|\alpha\rangle = \sqrt{\alpha_0} |00\rangle + \sqrt{\alpha_1} |11\rangle$

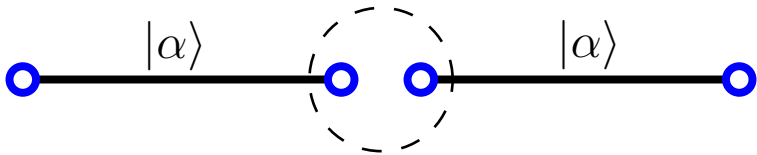
Local operations (and classical communication): qubits not allowed to interact

Maximally Entangled:  $|\Psi\rangle = \frac{1}{\sqrt{2}} |00\rangle + \frac{1}{\sqrt{2}} |11\rangle$

Singlet, Bell State, Maximally Entangled State **Singlet Conversion Probability**  $p = 2\alpha_1$ , for  $\alpha_0 > \alpha_1$

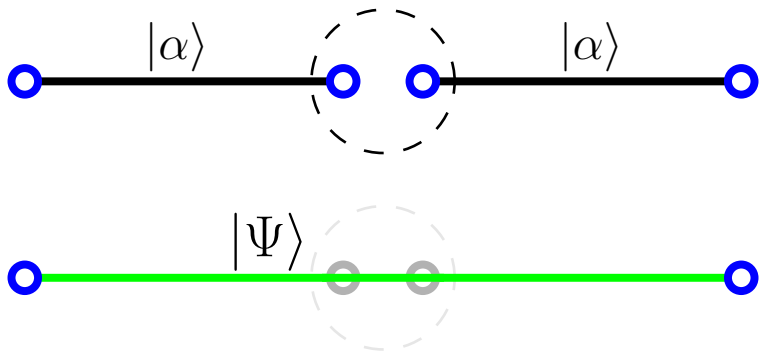
Otherwise: product state (failure)

# Entanglement Swapping



We can entangle the two outermost qubits, using only local operations and classical communication: *i.e.* without interacting outermost qubits. Using **entanglement swapping**.

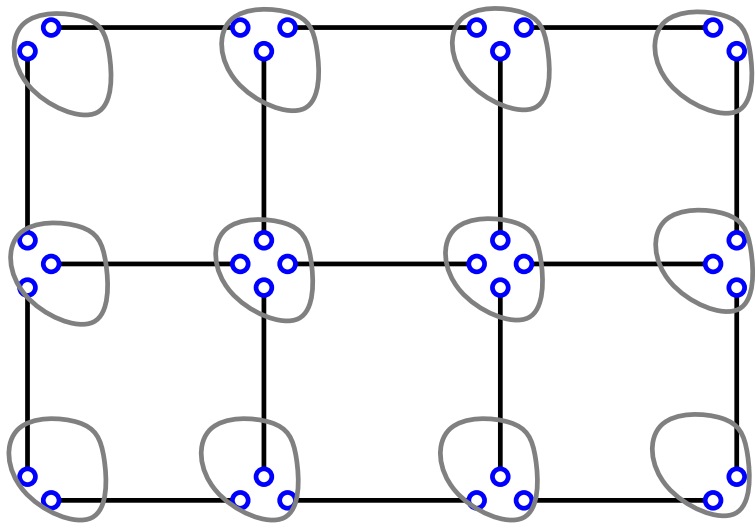
# Entanglement Swapping



Entanglement Swapping. Get **Bell state** with same probability as in singlet conversion  $p = 2\alpha_1$  ! (product state otherwise) Note: if  $\alpha_1 = 1/2$ , then  $p = 1$ .



# Quantum Network



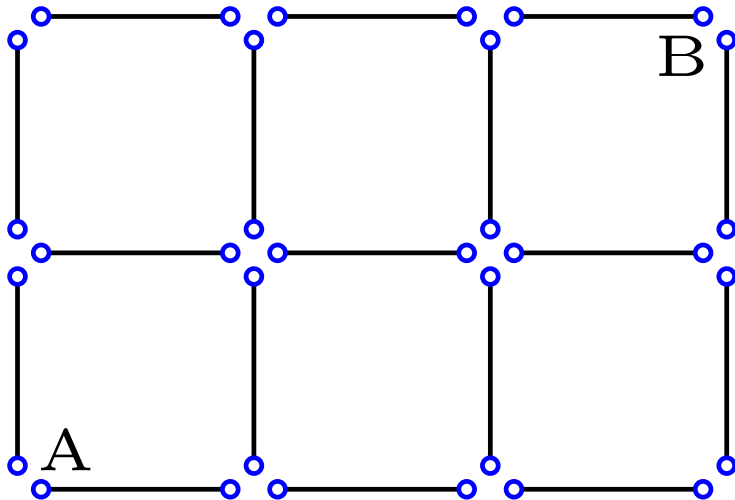
Concrete: Square lattice. Each bond is an entangled pair with amount of entanglement  $\alpha_1$ .

# Quantum Network

How to treat a network larger than two pairs. Most naive method: repeated swapping  $\Rightarrow$  exponential decay. Next most naive: borrow ideas from one-dimensional quantum repeaters.

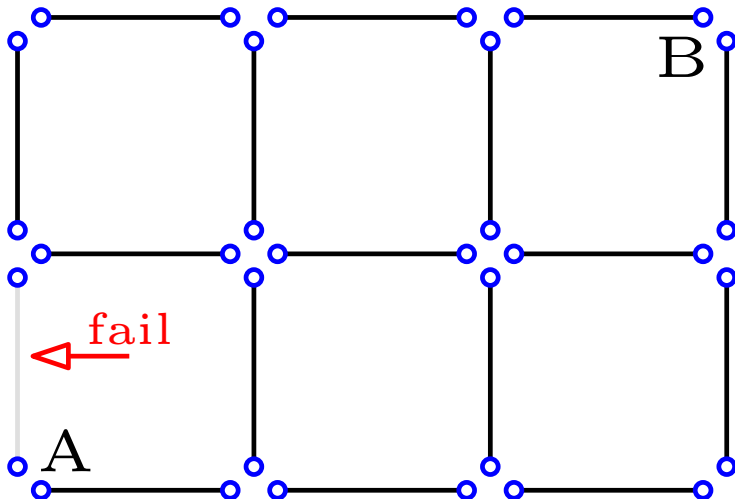
- 1 Attempt to put each pair in a Bell state. Here: Singlet conversion with probability of success  $p = 2\alpha_1$ .
- 2 Entanglement swappings between pairs of these Bell states. Result: New Bell state between outermost qubits, one from each of the pairs.
- 3 Repeat swappings, entangling ever more distant qubits.

# “Classical” Entanglement Percolation



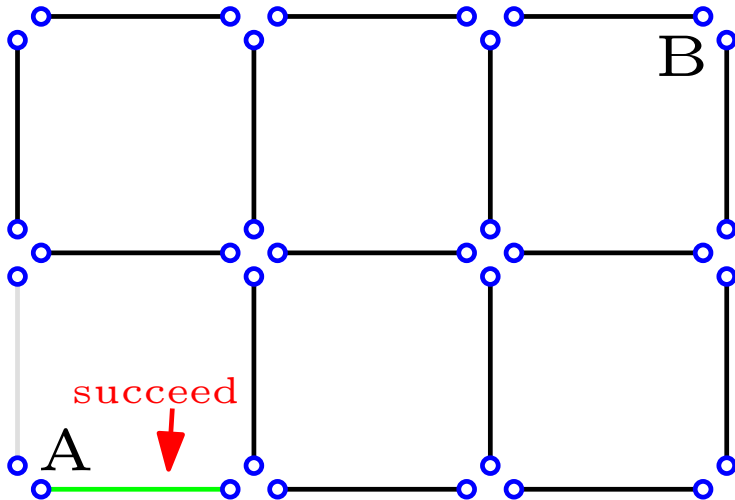
Entangle nodes A and B

# “Classical” Entanglement Percolation



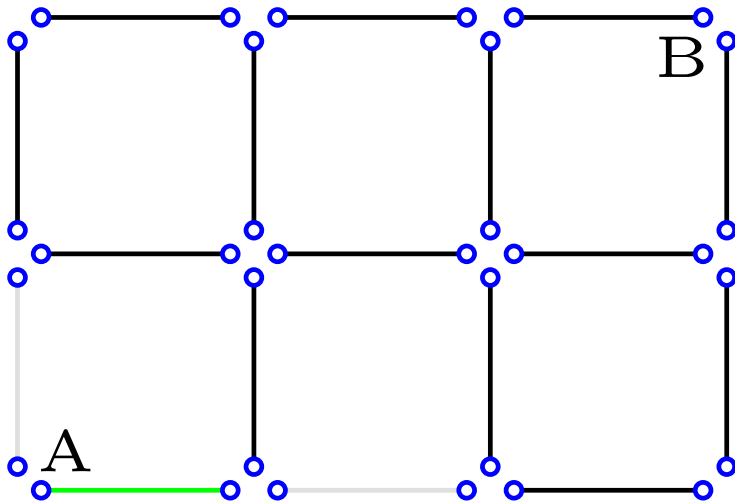
Singlet conversion fails.

# “Classical” Entanglement Percolation



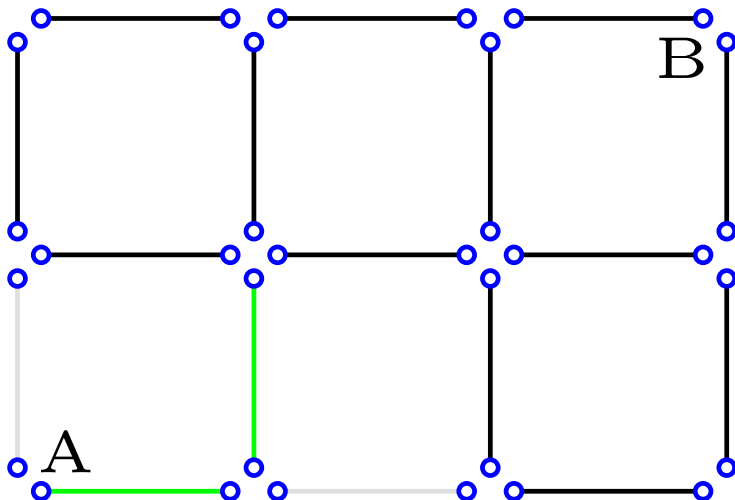
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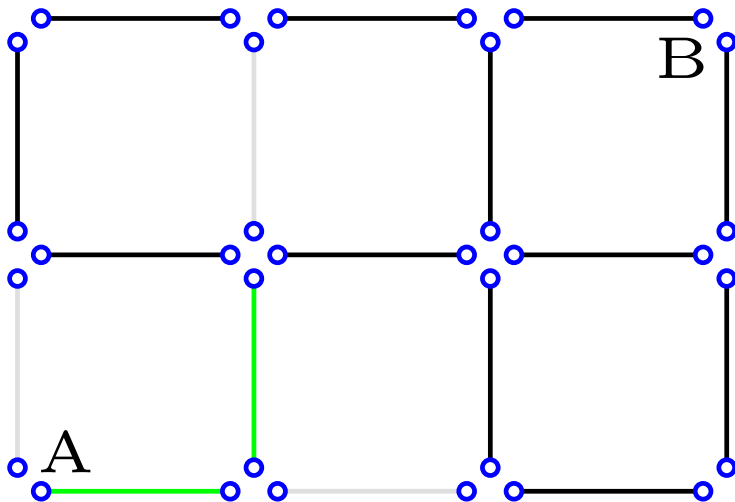
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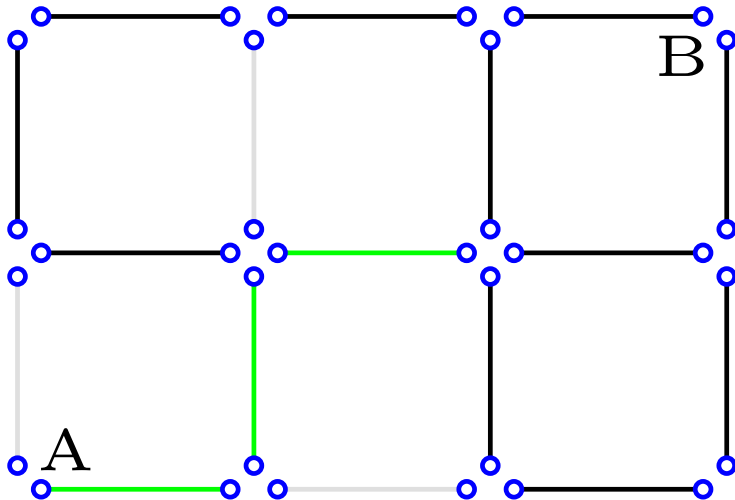
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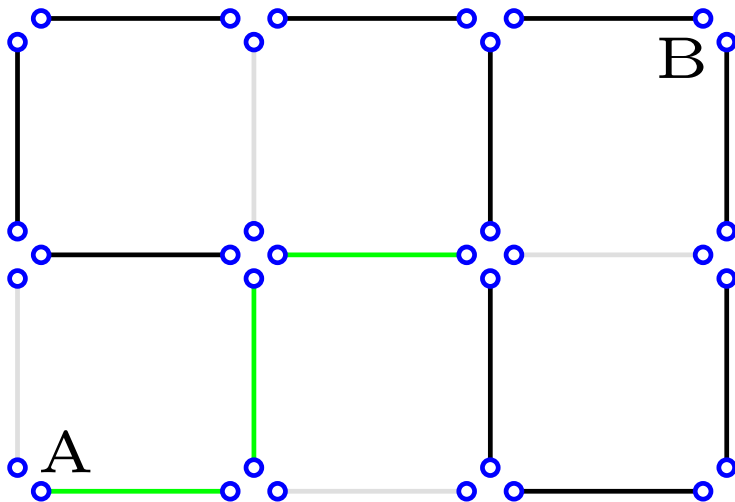


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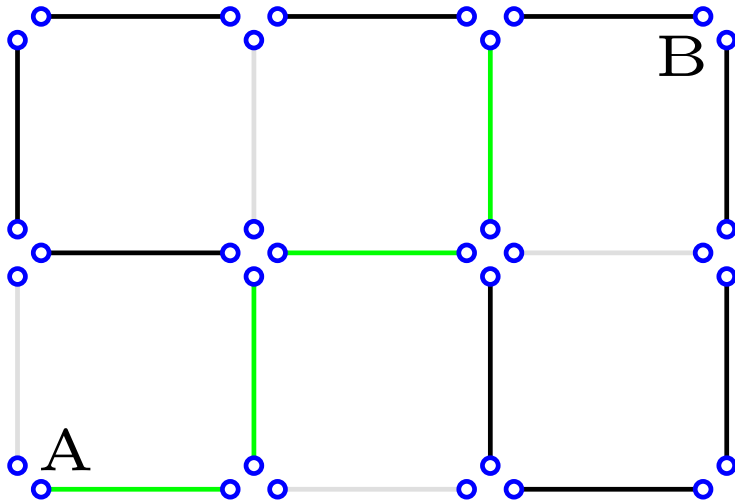
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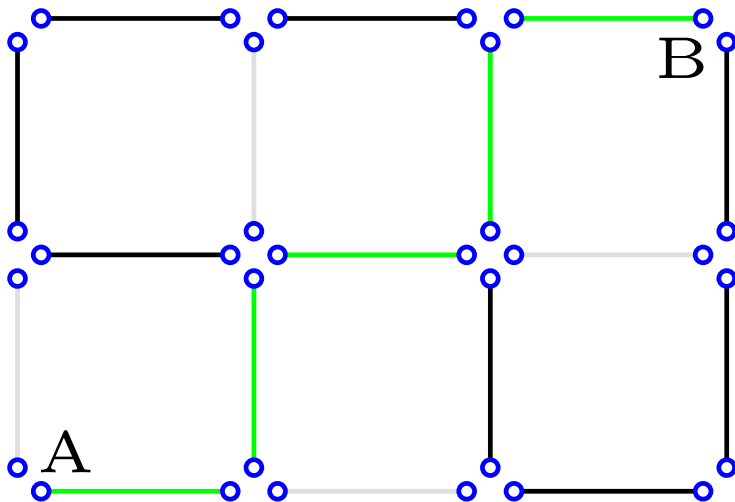
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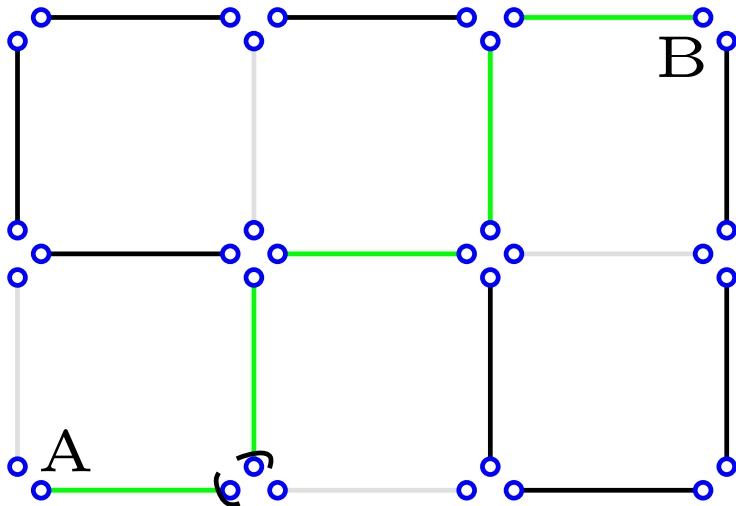
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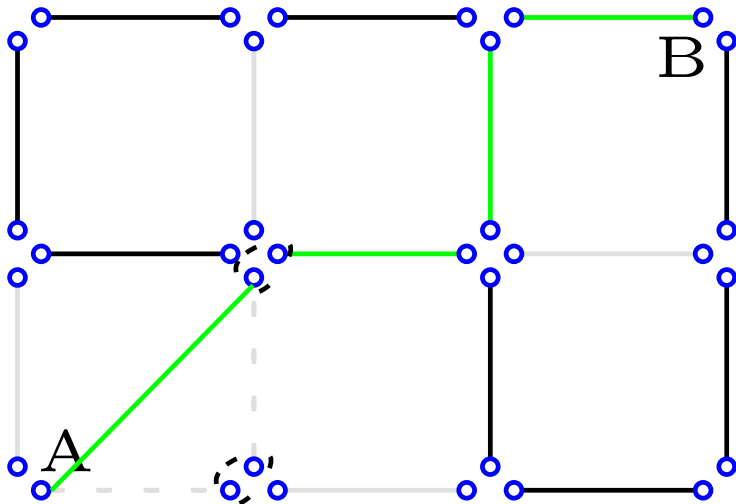
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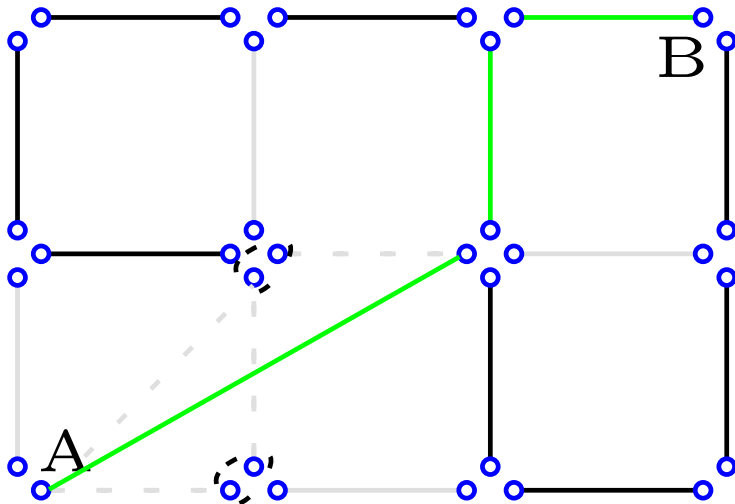
Entanglement swapping succeeds,  $p = 2\alpha_1 = 1$ .

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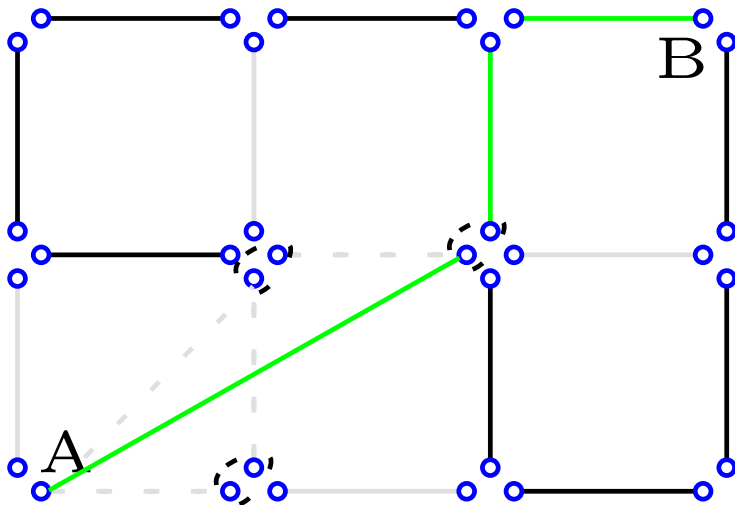
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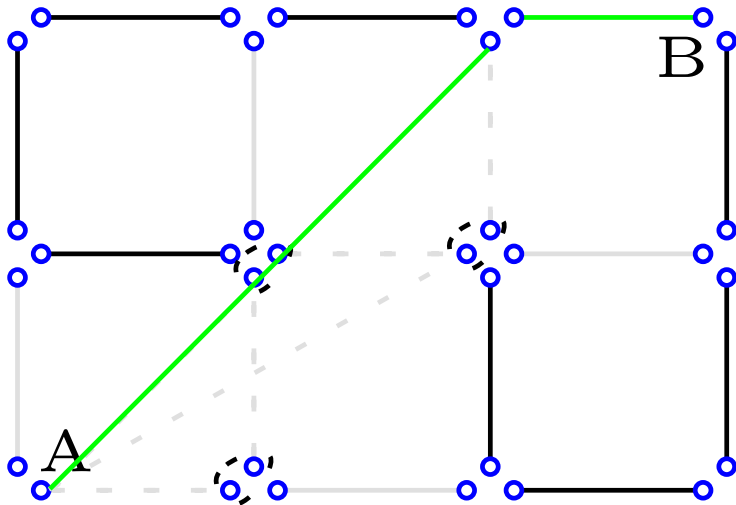
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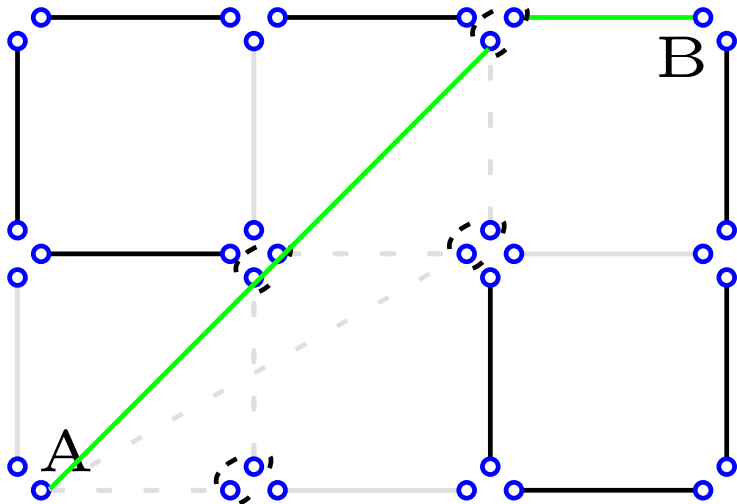


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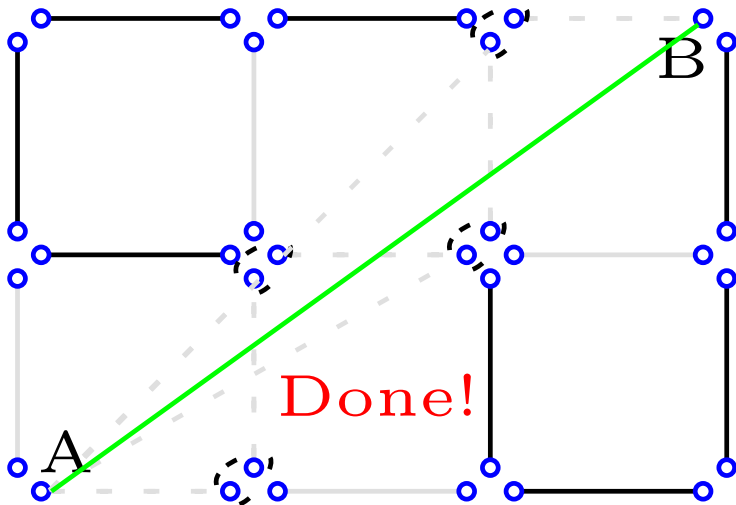
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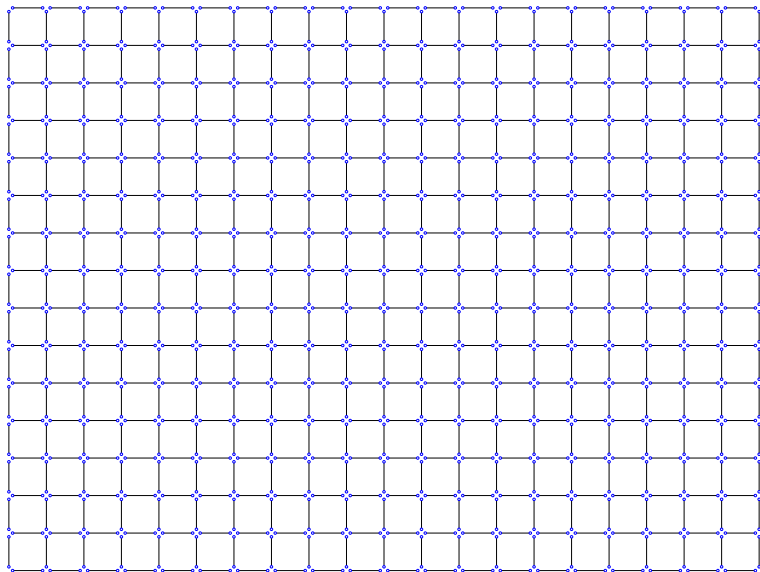
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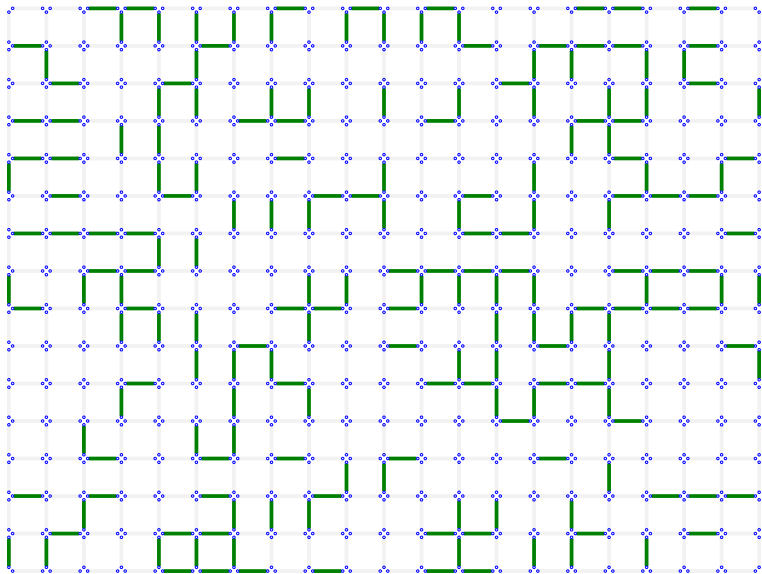
Entanglement swapping succeeds,  $p = 2\alpha_1 = 1$ .

Big Network:  $\alpha_1 = 0.175$   $p = 0.35$



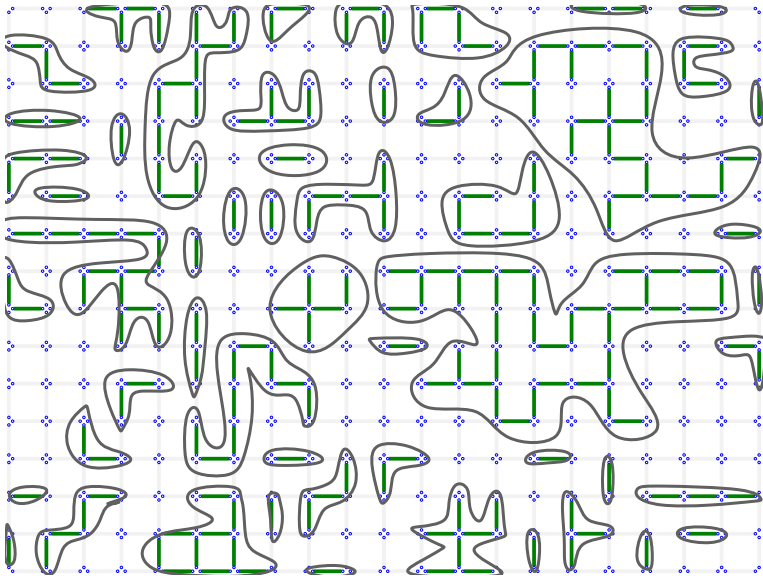
All bonds identically prepared in state  $|\alpha\rangle$ .

Big Network:  $\alpha_1 = 0.175$   $p = 0.35$



Singlet conversions everywhere... Partition into clusters.

Big Network:  $\alpha_1 = 0.175$   $p = 0.35$

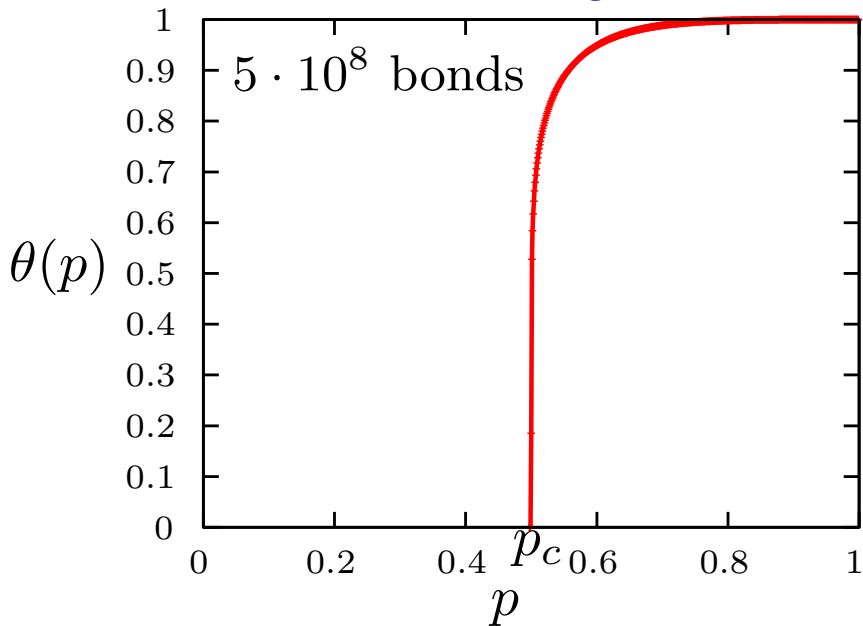


Bond percolation on square lattice. Infinite cluster iff  $p > p_c = 0.5$ .

# Percolation theory

- Bonds are *open* (present) with probability  $p$ ; or else *closed* (absent).  $p$  is called the bond density.
- For large lattices there is a threshold value of the bond density  $p_c$ . For  $p > p_c$  there is a single cluster that spans the whole lattice.  $p_c$  depends on the structure of the lattice.
- In order for A and B to be connected if they are very far apart,
  - ① Must have  $p > p_c$
  - ② Both A and B must be in the huge cluster
- Let  $\theta(p) = \text{Prob}(A \text{ is in huge cluster})$
- A and B are connected with  $\text{Prob} = \theta^2(p)$

# Fraction of Bonds in largest cluster





Can we do better than simply swapping along a chain ?

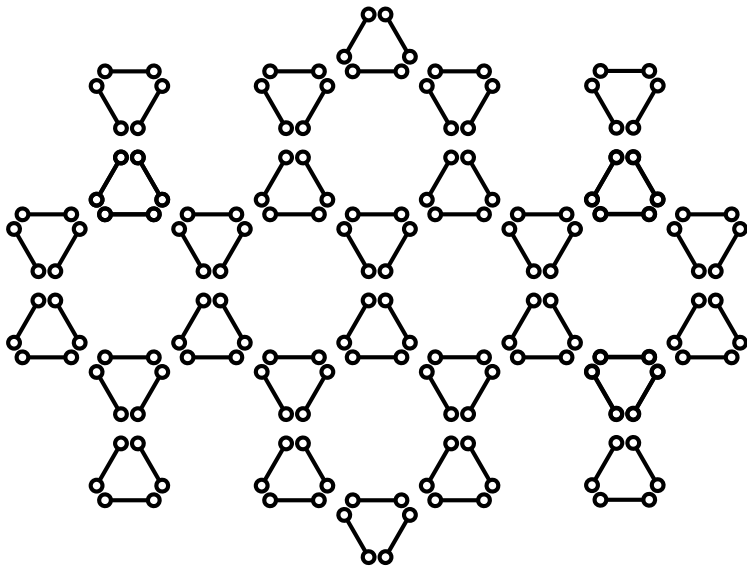
Can we do better than simply swapping along a chain ?

Yes. Precondition the lattice with other quantum operations.

Change local structure  $\Rightarrow$  Different lattice  $\Rightarrow$  Different global properties: Different percolation threshold.

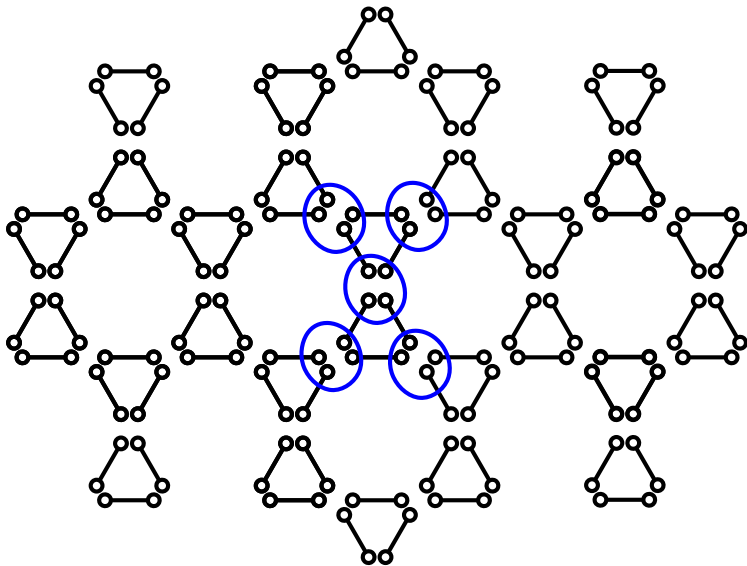
Then swap along chain.

# Kagome lattice



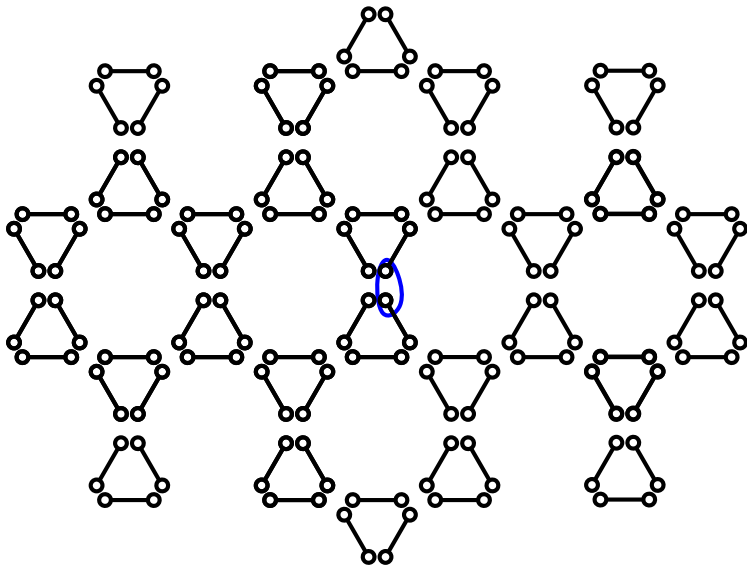
Easiest example.

# Kagome lattice



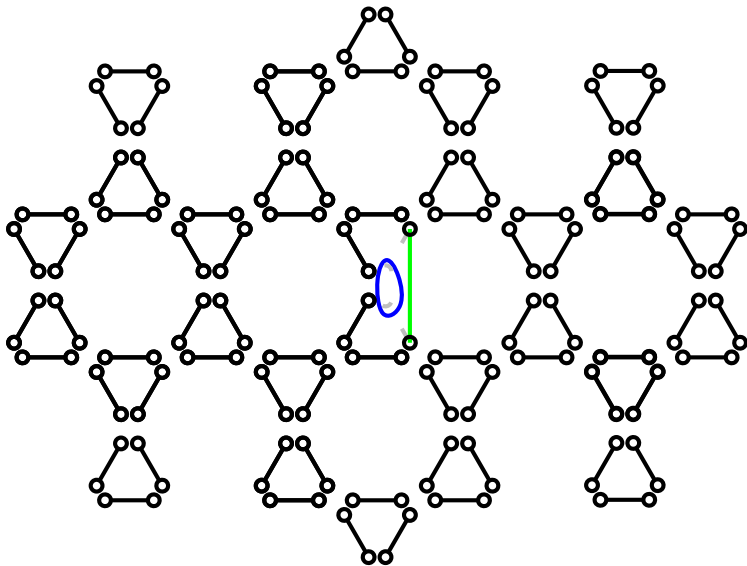
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## Kagome lattice



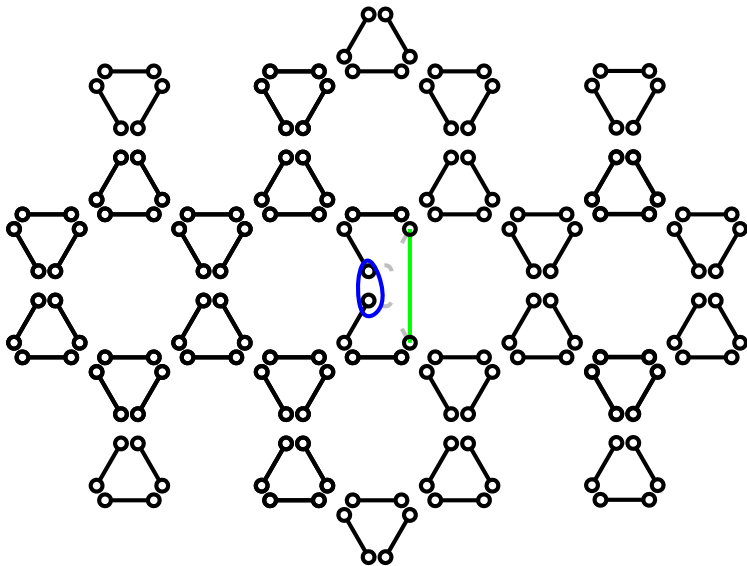
Entanglement swapping to create vertical bonds.

## Kagome lattice



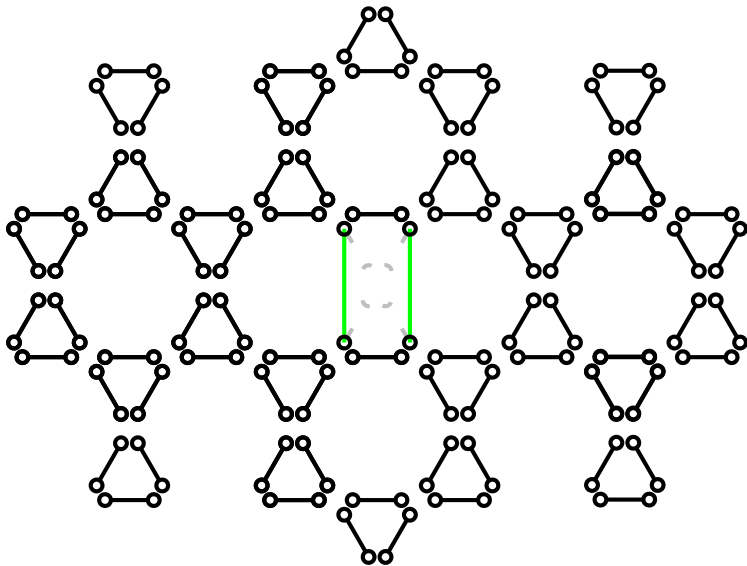
Vertical bonds are Bell pairs.

## Kagome lattice



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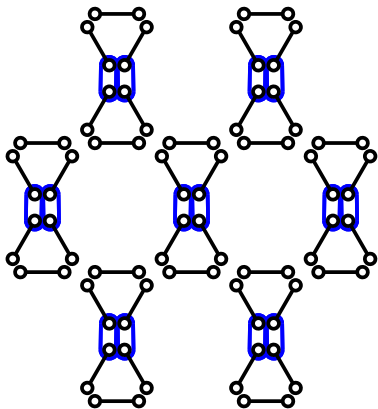
## Kagome lattice



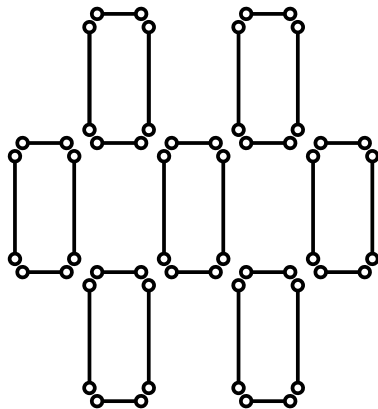
Then singlet conversion on horizontal bonds.



## Convert kagome lattice to square lattice



$p_c \approx 0.52$  Kagome



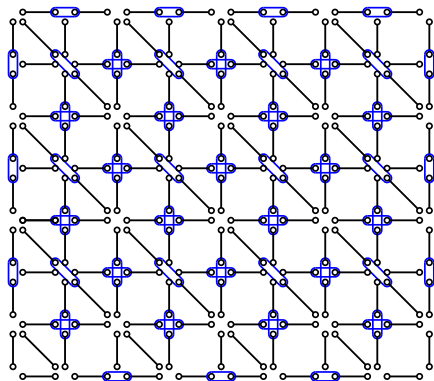
$p_c = 0.5$  Square lattice

*Acín, Cirac, Lewenstein, Nature Phys (2007)*

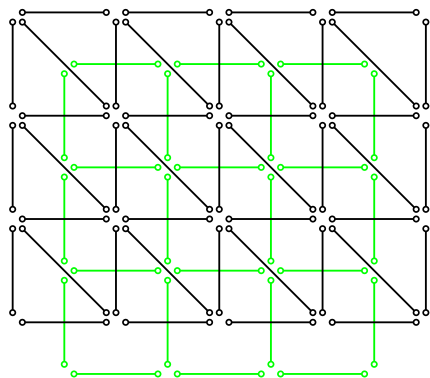
*Perseguers, Cirac, Acín, Lewenstein, Wehr, PRA (2008)*

*JL, Wehr, Lewenstein, PRA (2009)*

## Convert bowtie lattice to square and triangular



$p_c \approx 0.40$  Bowtie



$p_c \approx 0.35$  Triangular

*JL, J. Wehr, M. Lewenstein, PRA (2009)*

# Characterize Protocols

In all known effective preconditioning protocols:

- Local connectivity non-decreasing: Coordination number increases or remains the same.
- Global connectivity non-decreasing: Classical percolation threshold decreases.

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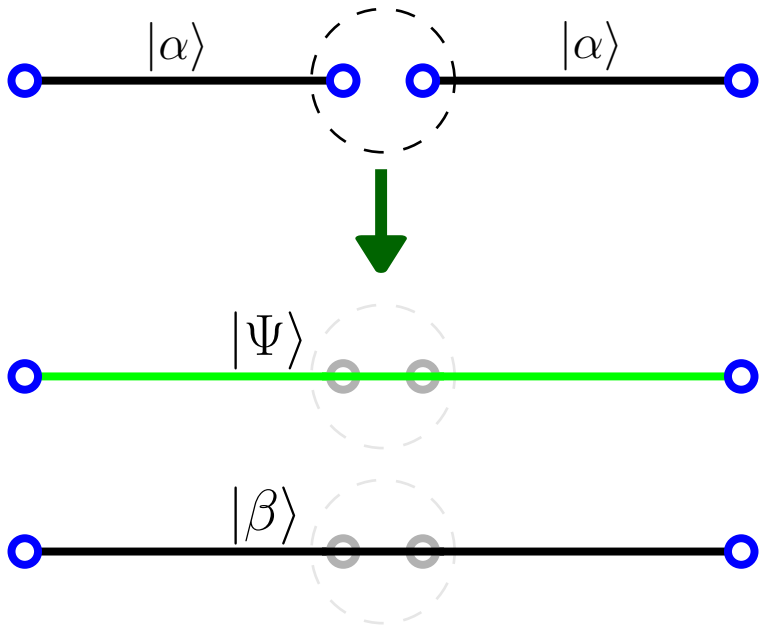
**No.** Counter-example: Look at swapping more closely.  
Only perform partial entanglement swapping.

# Entanglement Swapping

Entanglement swapping procedure.

- 1 Measurement projecting two center qubits onto Bell basis.
- 2 Unitary on end qubit based on result of measurement, leaving two distinct (pairs of) states.
- 3 One state is partially entangled: perform SCP on it.
- 4 Other state is already maximally entangled.
- 5 Averaging over these two possibilities gives  $p = 2\alpha_1$ .

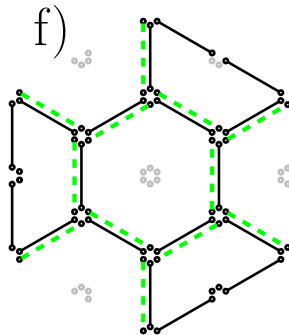
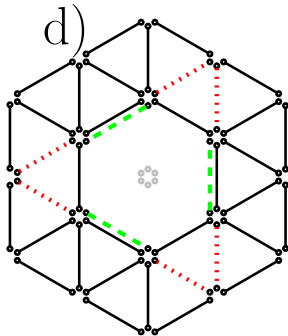
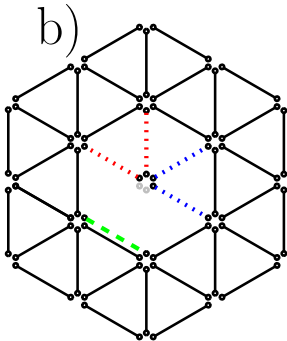
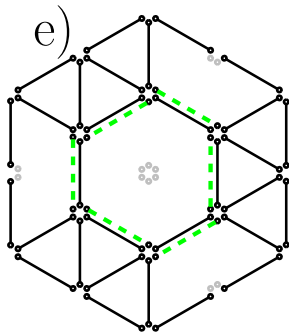
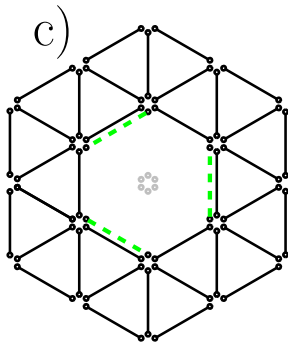
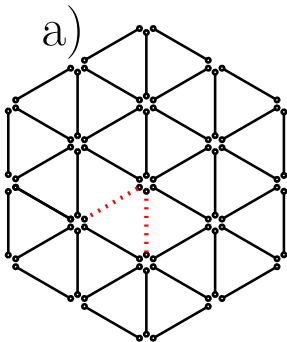
# Entanglement Swapping



# Partial swapping and distillation

Do partial swapping (only projection) on selected bonds of triangular lattice. This leaves pairs of parallel bonds in honeycomb (hexagonal) lattice, which are then distilled.





# Partial swapping and distillation

Majorization says we can distill a singlet from the double bond pair with probability

$$p = \min \left\{ 1, 2 \left( 1 - \frac{\alpha_0^3}{\alpha_0^2 + \alpha_1^2} \right) \right\}$$

The real root of  $\alpha_0^3 - \alpha_0^2 + \alpha_0 - 1/2 = 0$  is  $\alpha_0^* \approx 0.647798$ . Thus, if  $\alpha_0 < 0.647$  then each double bond can be converted to a singlet with probability 1.

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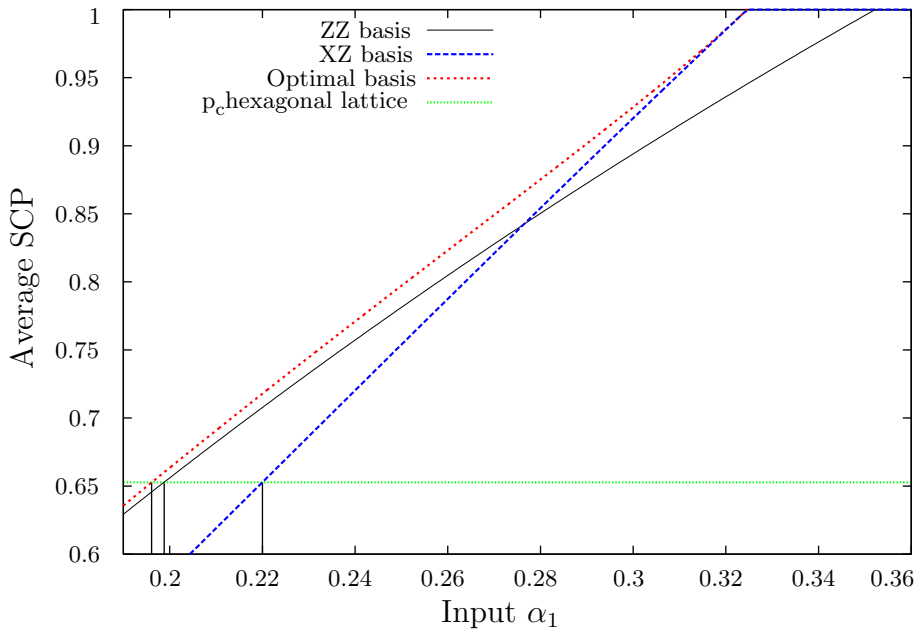
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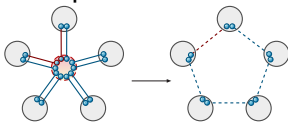
There is a range of initial entanglement for which long-range entanglement is achieved with probability 1.

And this can be improved by optimizing the chosen Bell basis instead of taking the usual one for swapping.



# Other work

- Multipartite (GHZ) states  $\Rightarrow$  percolation non-planar graphs  
*Perseguers, Cavalcanti, JL, Lewenstein, and Acín, PRA (2010)*
- Mixed states of rank  $\leq 3$  *Broadfoot, Dorner, Jaksch, PRA (2010), EPL (2009)*
- Q-star transformation on complex networks. increase



entanglement distance.

*Cuquet,*

- Calsamiglia, PRL (2009), PRA (2011)*
- Mixed states full rank, complex network. *JL, Perseguers, Lewenstein, Acín, QIC (2012)*
- Review: *Perseguers, JL, Cavalcanti, Lewenstein, Acín, Rep. Prog. Phys. 76 (2013)*

## Other things ...

- Geometry, topology, etc. often determines optimal protocol.
- Constraint on geometry of transformed lattice?
- Is there a lower bound on local entanglement density in 2-d ?
- Dynamics ?
- Coupled lattices ?